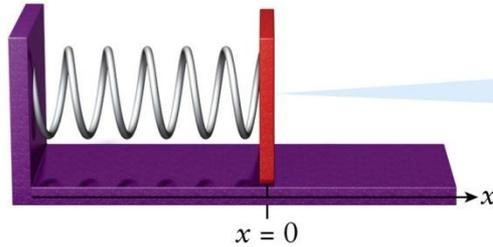


# Chapter 12

## Oscillations

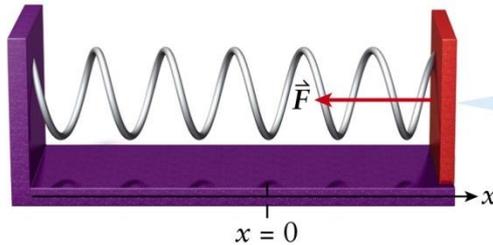


(a)



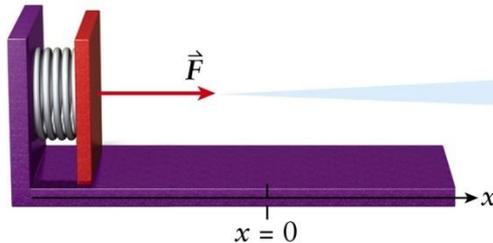
- ① The object at the free end of the spring is in equilibrium when the spring is neither stretched nor compressed. We call this point  $x = 0$ .

(b)



- ② When the spring is stretched, the object is at  $x > 0$ . The spring force pulls the object back toward equilibrium ( $F_x < 0$ ).

(c)



- ③ When the spring is compressed, the object is at  $x < 0$ . The spring force pushes the object back toward equilibrium ( $F_x > 0$ ).



The force exerted by the spring is a restoring force: No matter which way the object is displaced from equilibrium, the spring force always acts to return the object to equilibrium.

Period  $T$  of an oscillation

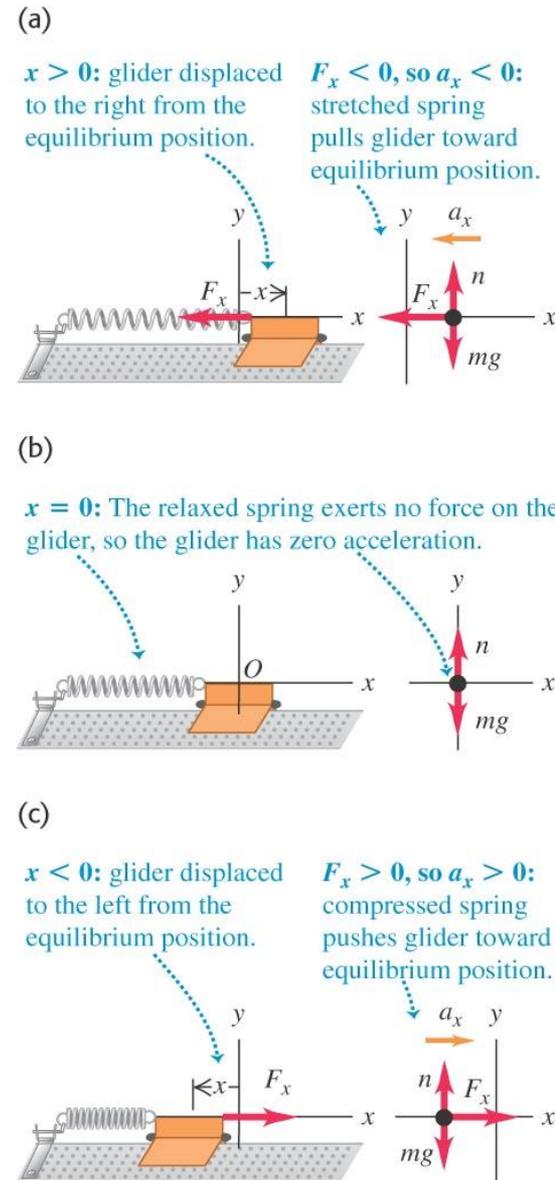
$$f = \frac{1}{T} \quad \text{and} \quad T = \frac{1}{f}$$

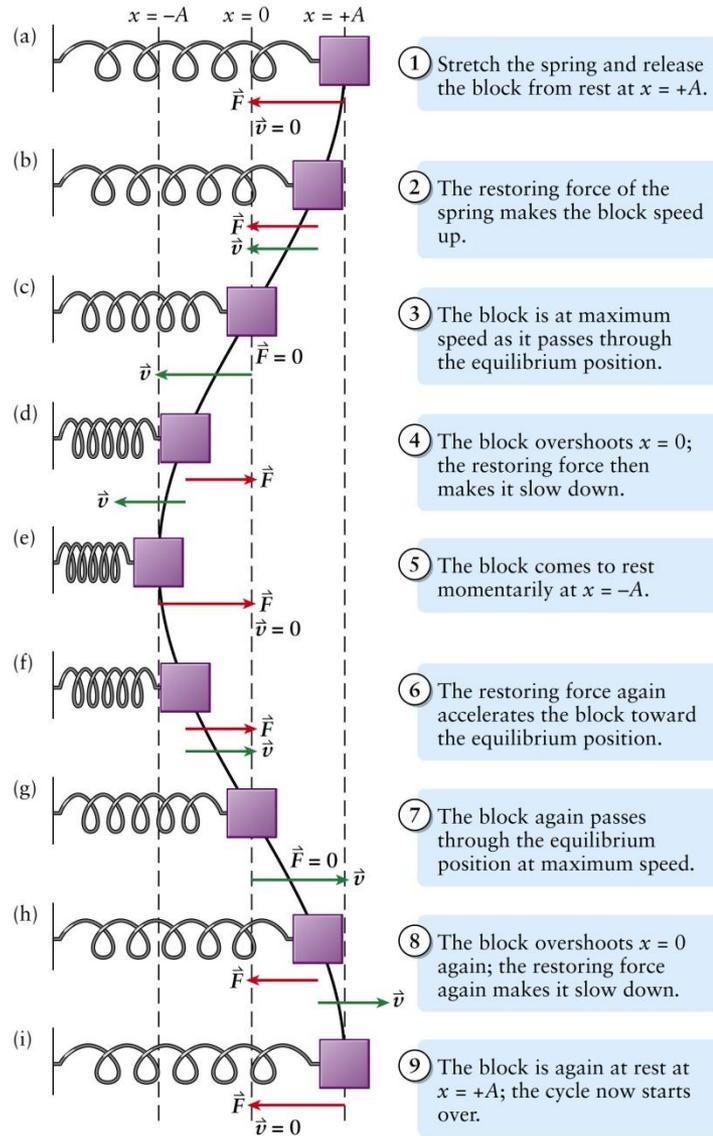
Frequency  $f$  of the oscillation

- $fT=1$  – freq  $f$ (Hz) time period  $T$ (s) =1
- **$f=1/T$**     $\omega = 2\pi f$     $\omega T=2\pi$     **$T=2\pi / \omega$**

# What causes periodic motion?

- If a body attached to a spring is displaced from its equilibrium position, the spring exerts a *restoring force* on it, which tends to restore the object to the equilibrium position. This force causes *oscillation* of the system, or *periodic motion*.
- Figure at the right illustrates the restoring force  $F_x$ .





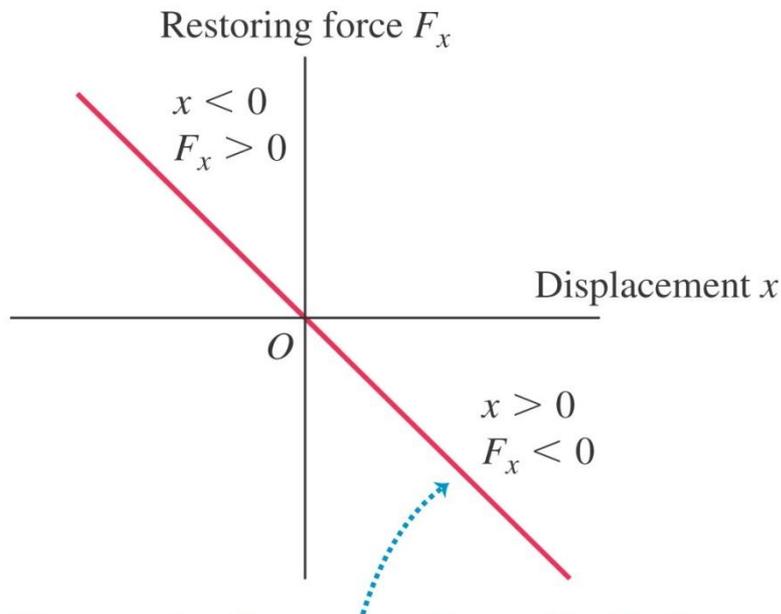
# Characteristics of periodic motion

- The *amplitude*,  $A$ , is the maximum magnitude of displacement from equilibrium.
- The *period*,  $T$ , is the time for one cycle.
- The *frequency*,  $f$ , is the number of cycles per unit time.
- The *angular frequency*,  $\omega$ , is  $2\pi$  times the frequency:  $\omega = 2\pi f$ .
- The frequency and period are reciprocals of each other:  
 $f = 1/T$  and  $T = 1/f$ .

# Simple harmonic motion (SHM)

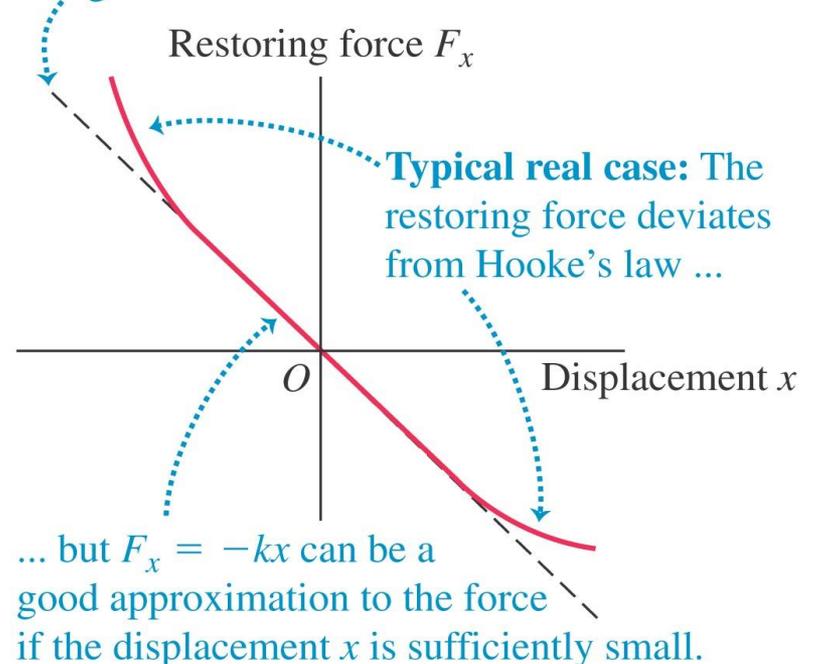
## Simple Harmonic Oscillator (SHO)

- When the restoring force is *directly proportional* to the displacement from equilibrium, the resulting motion is called *simple harmonic motion* (SHM).
- An ideal spring obeys Hooke's law, so the restoring force is  $F_x = -kx$ , which results in simple harmonic motion.



The restoring force exerted by an idealized spring is directly proportional to the displacement (Hooke's law,  $F_x = -kx$ ): the graph of  $F_x$  versus  $x$  is a straight line.

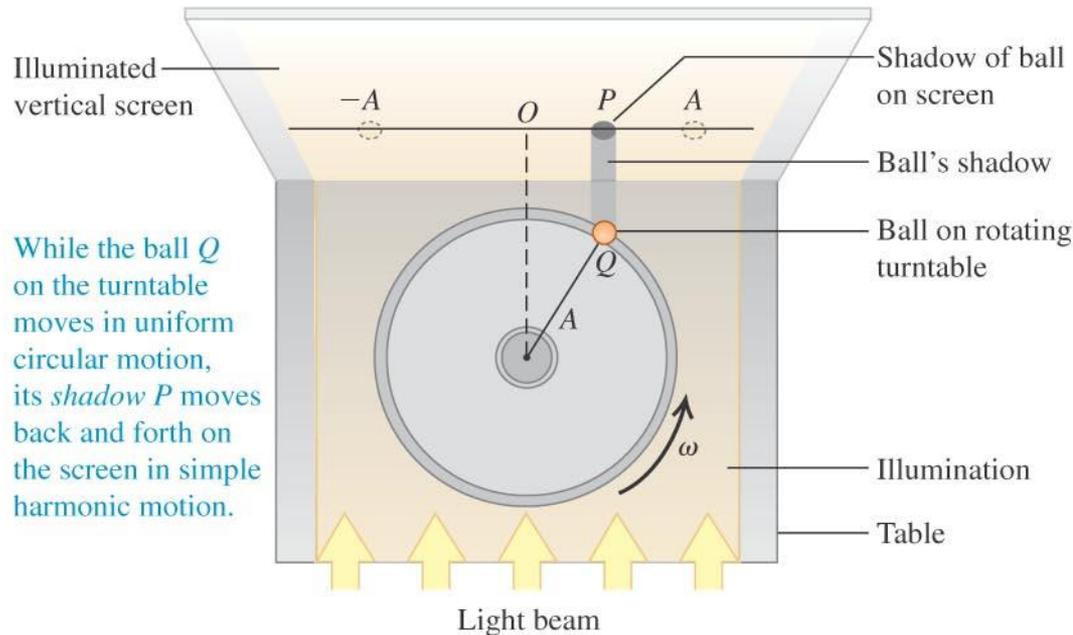
**Ideal case:** The restoring force obeys Hooke's law ( $F_x = -kx$ ), so the graph of  $F_x$  versus  $x$  is a straight line.



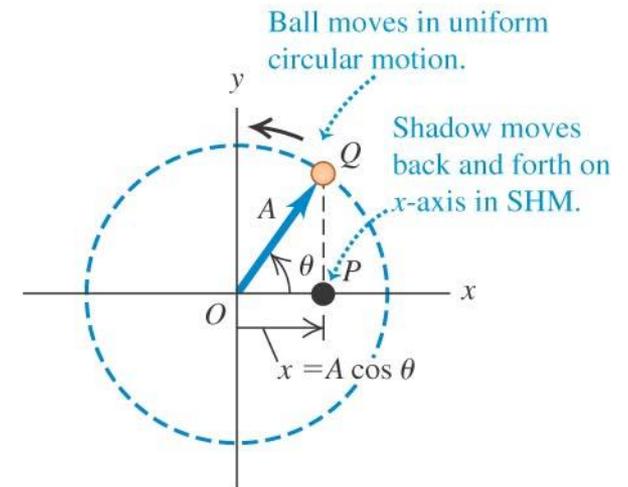
# Simple harmonic motion viewed as a projection

- Simple harmonic motion is the projection of uniform circular motion onto a diameter

(a) Apparatus for creating the reference circle



(b) An abstract representation of the motion in (a)

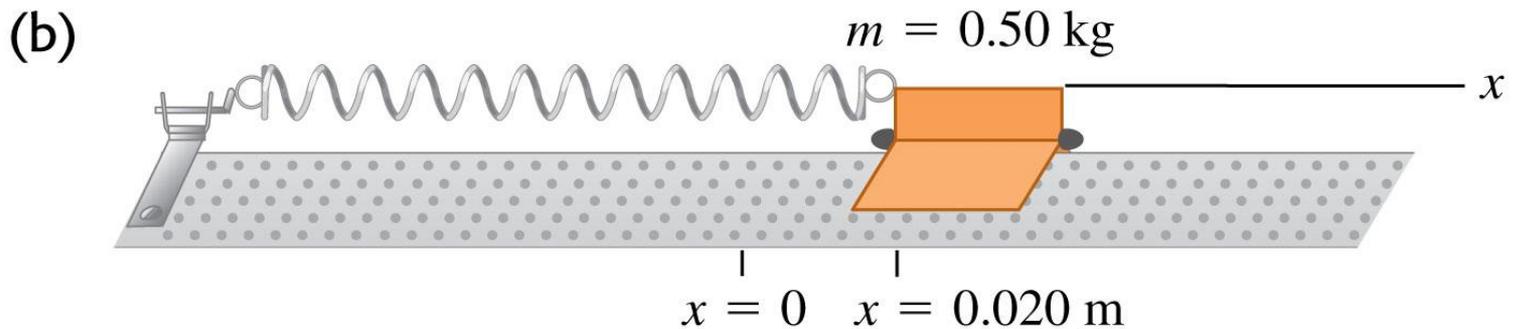
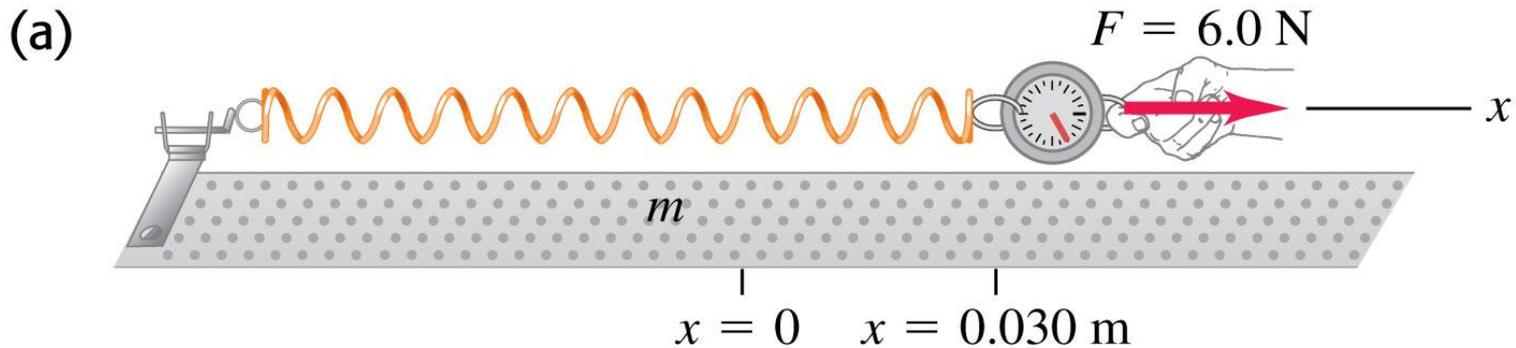


# Characteristics of SHM

- For a body vibrating by an ideal spring:

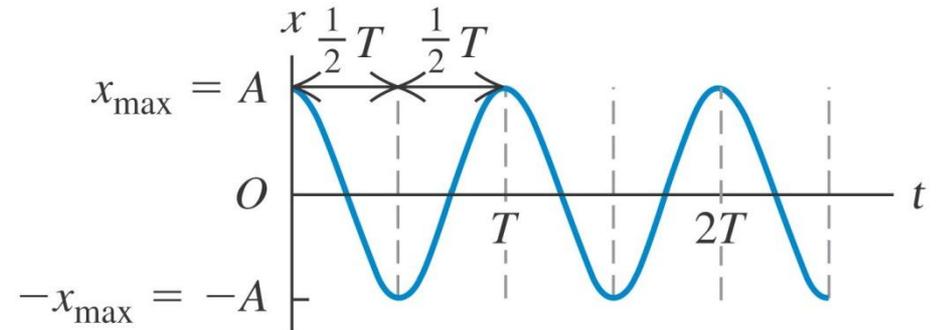
$$\omega = \sqrt{\frac{k}{m}} \quad f = \frac{\omega}{2\pi} = \frac{1}{2\pi} \sqrt{\frac{k}{m}} \quad T = \frac{1}{f} = \frac{2\pi}{\omega} = 2\pi \sqrt{\frac{m}{k}}$$

- Follow Example 14.2 and Figure 14.8 below.



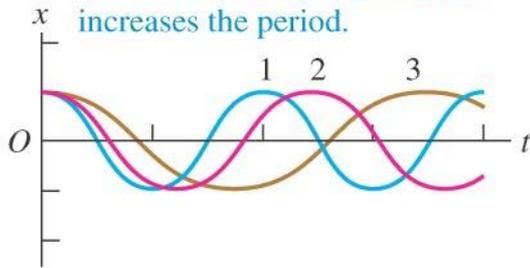
# Displacement as a function of time in SHM

- The displacement as a function of time for SHM with phase angle  $\phi$  is  $x = A\cos(\omega t + \phi)$
- Changing  $m$ ,  $A$ , or  $k$  changes the graph of  $x$  versus  $t$ , as shown below.



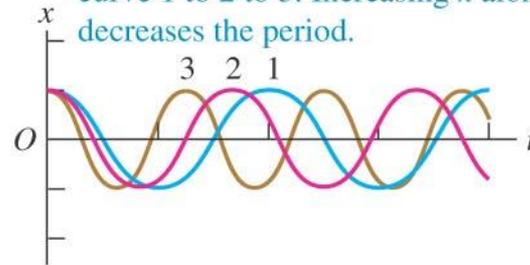
(a) Increasing  $m$ ; same  $A$  and  $k$

Mass  $m$  increases from curve 1 to 2 to 3. Increasing  $m$  alone increases the period.



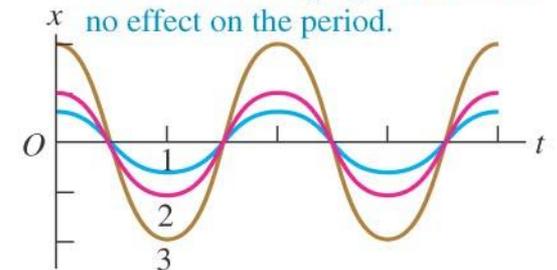
(b) Increasing  $k$ ; same  $A$  and  $m$

Force constant  $k$  increases from curve 1 to 2 to 3. Increasing  $k$  alone decreases the period.



(c) Increasing  $A$ ; same  $k$  and  $m$

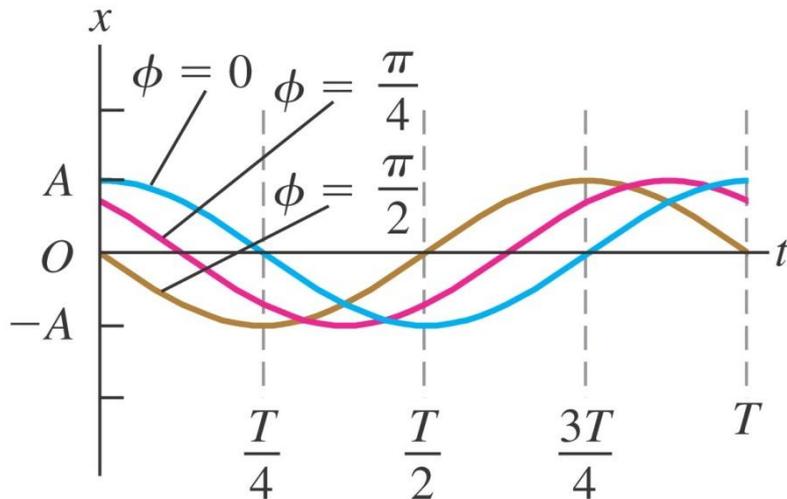
Amplitude  $A$  increases from curve 1 to 2 to 3. Changing  $A$  alone has no effect on the period.



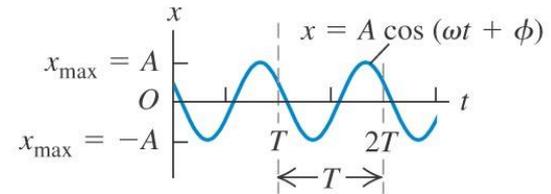
# Displacement, velocity, and acceleration

- The graphs below show  $x$ ,  $v_x$ , and  $a_x$  for  $\phi = \pi/3$ .
- The graph below shows the effect of different phase angles.

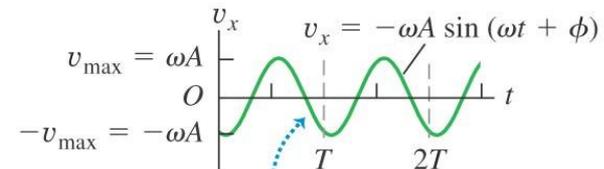
These three curves show SHM with the same period  $T$  and amplitude  $A$  but with different phase angles  $\phi$ .



(a) Displacement  $x$  as a function of time  $t$

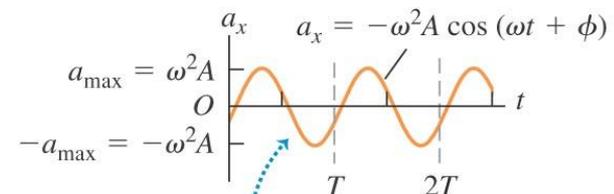


(b) Velocity  $v_x$  as a function of time  $t$



The  $v_x-t$  graph is shifted by  $\frac{1}{4}$  cycle from the  $x-t$  graph.

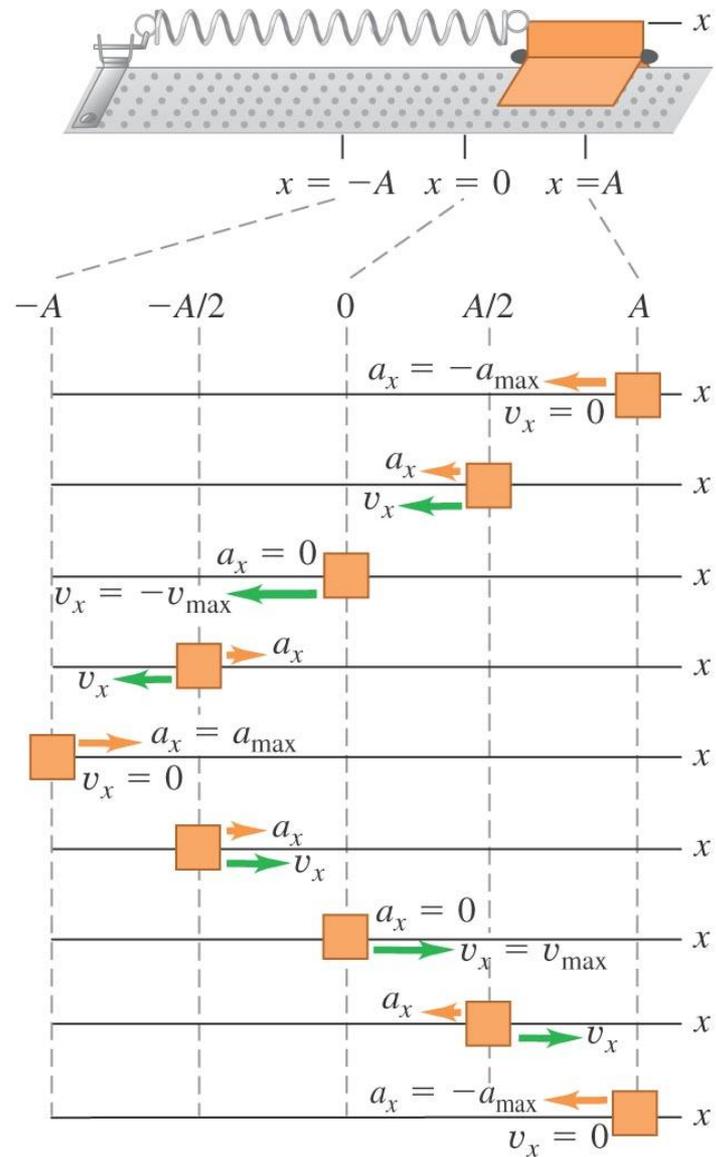
(c) Acceleration  $a_x$  as a function of time  $t$



The  $a_x-t$  graph is shifted by  $\frac{1}{4}$  cycle from the  $v_x-t$  graph and by  $\frac{1}{2}$  cycle from the  $x-t$  graph.

# Behavior of $v_x$ and $a_x$ during one cycle

- Figure shows how  $v_x$  and  $a_x$  vary during one cycle.



## SHO - mass and amplitude

An object on the end of a spring is oscillating in simple harmonic motion. If the amplitude of oscillation is doubled, how does this affect the oscillation period  $T$  and the object's maximum speed  $v_{\max}$ ?

- A.  $T$  and  $v_{\max}$  both double.
- B.  $T$  remains the same and  $v_{\max}$  doubles.
- C.  $T$  and  $v_{\max}$  both remain the same.
- D.  $T$  doubles and  $v_{\max}$  remains the same.
- E.  $T$  remains the same and  $v_{\max}$  increases by a factor of

$$\sqrt{2}.$$

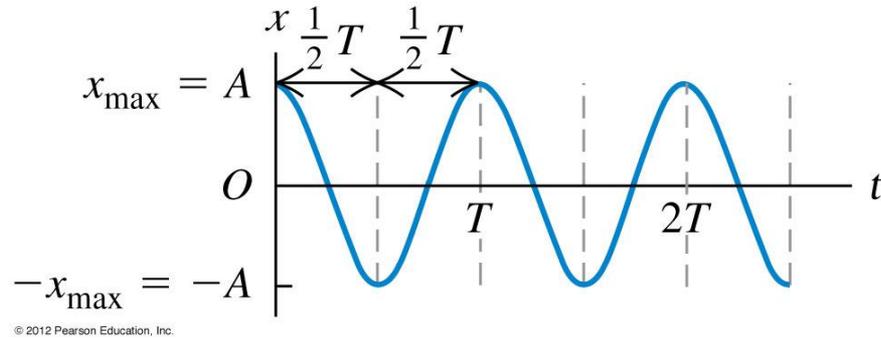
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- E.  $T$  remains the same and  $v_{\max}$  increases by a factor of  $\sqrt{2}$ .

$$\sqrt{2}$$

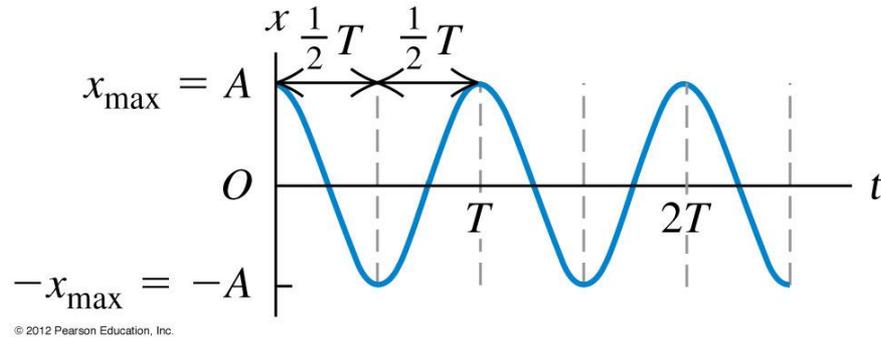
This is an  $x-t$  graph for an object in simple harmonic motion.



At which of the following times does the object have the *most negative velocity*  $v_x$ ?

- A.  $t = T/4$
- B.  $t = T/2$
- C.  $t = 3T/4$
- D.  $t = T$
- E. Two of the above are tied for most negative velocity

This is an  $x$ - $t$  graph for an object in simple harmonic motion.



At which of the following times does the object have the *most negative velocity*  $v_x$ ?



A.  $t = T/4$

B.  $t = T/2$

C.  $t = 3T/4$

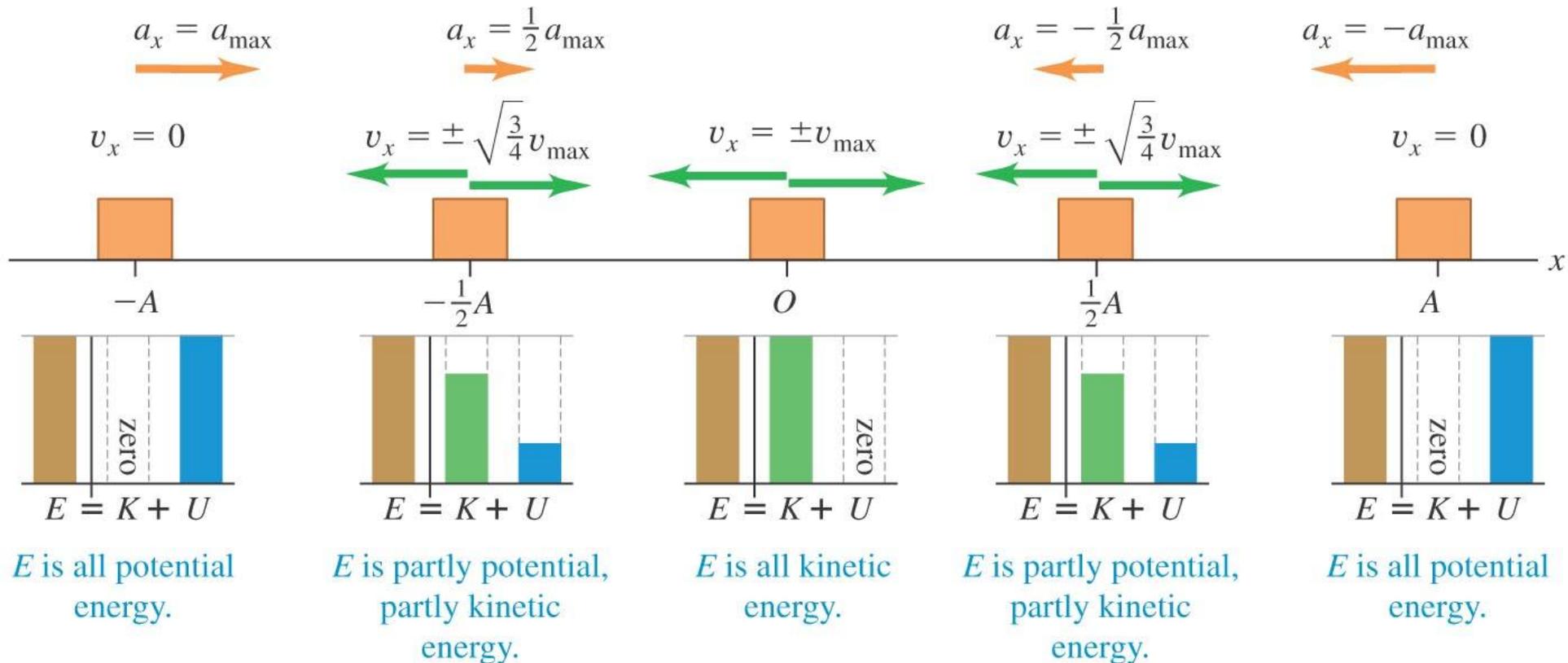
D.  $t = T$

E. Two of the above are tied for most negative velocity

# Energy in SHM

- The total mechanical energy  $E = K + U$  is conserved in SHM:

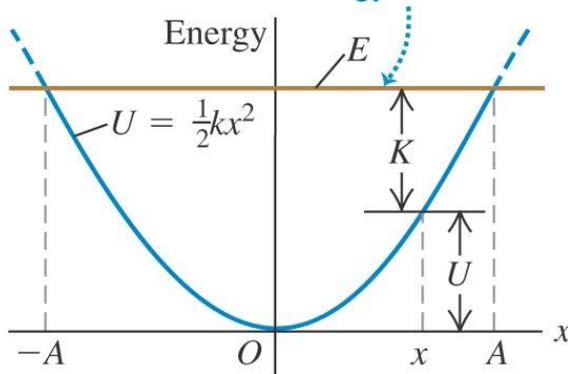
$$E = \frac{1}{2} m v_x^2 + \frac{1}{2} k x^2 = \frac{1}{2} k A^2 = \frac{1}{2} m v_{x\text{-maximum}}^2 = \text{constant}$$



# Energy diagrams for SHM

(a) The potential energy  $U$  and total mechanical energy  $E$  for a body in SHM as a function of displacement  $x$

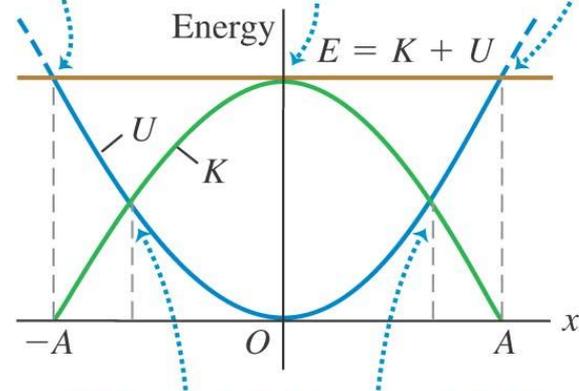
The total mechanical energy  $E$  is constant.



(b) The same graph as in (a), showing kinetic energy  $K$  as well

At  $x = \pm A$  the energy is all potential; the kinetic energy is zero.

At  $x = 0$  the energy is all kinetic; the potential energy is zero.



At these points the energy is half kinetic and half potential.

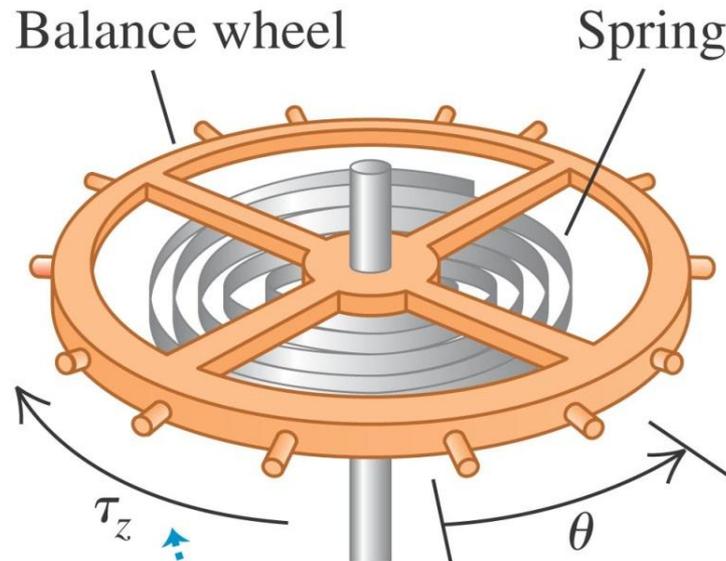
# Vertical SHM – Mass and Spring

## Gravity does NOT matter here

- If a body oscillates vertically from a spring, the restoring force has magnitude  $kx$ . Therefore the vertical motion is SHM.
- For a pendulum Gravity DOES matter.

# Angular SHM – old mechanical watch

- A coil spring exerts a restoring torque  $\tau_z = -\kappa\theta$ , where  $\kappa$  is called the *torsion constant* of the spring.
- The result is *angular* simple harmonic motion.



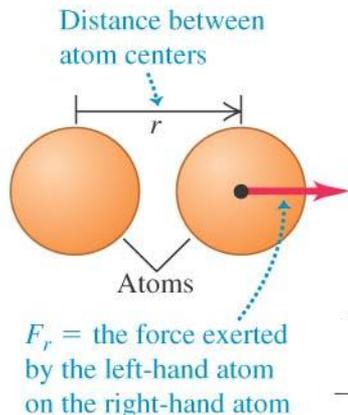
The spring torque  $\tau_z$  opposes the angular displacement  $\theta$ .

# Vibrations of molecules

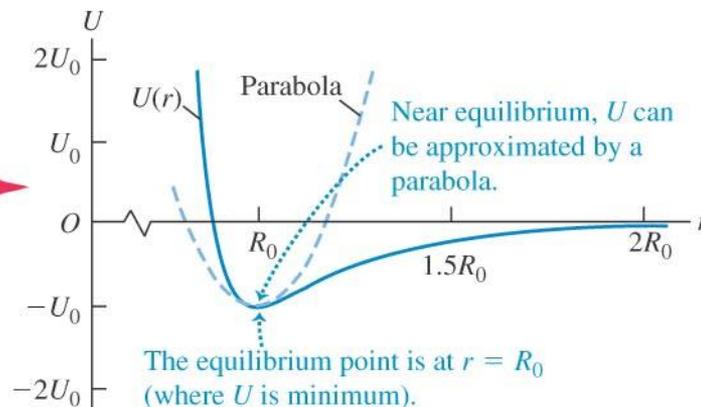
## Intermolecular forces

- Figure shows two atoms having centers a distance  $r$  apart, with the equilibrium point at  $r = R_0$ .
- If they are displaced a small distance  $x$  from equilibrium, the restoring force is  $F_r = -(72U_0/R_0^2)x$ , so  $k = 72U_0/R_0^2$  and the motion is SHM.
- Van der Waal like forces.

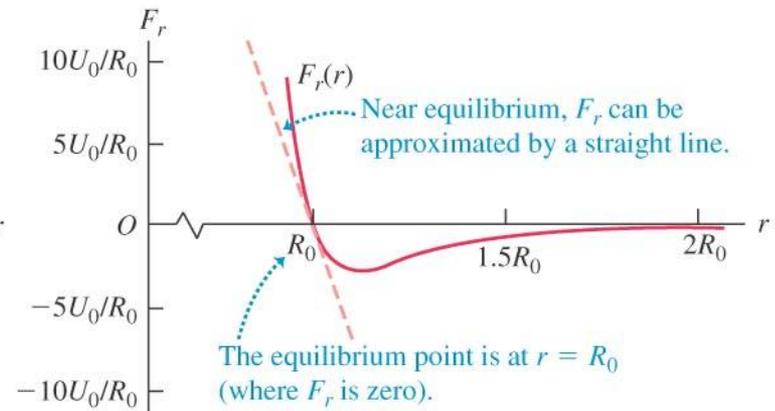
(a) Two-atom system



(b) Potential energy  $U$  of the two-atom system as a function of  $r$



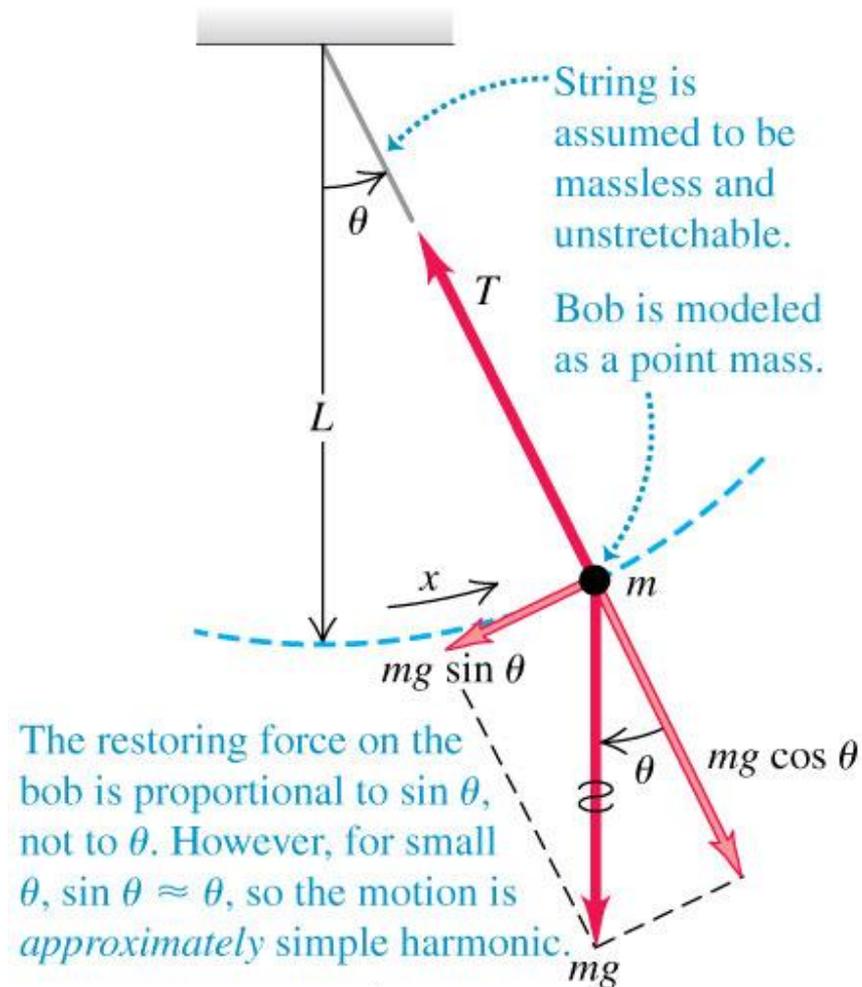
(c) The force  $F_r$  as a function of  $r$



# The simple pendulum

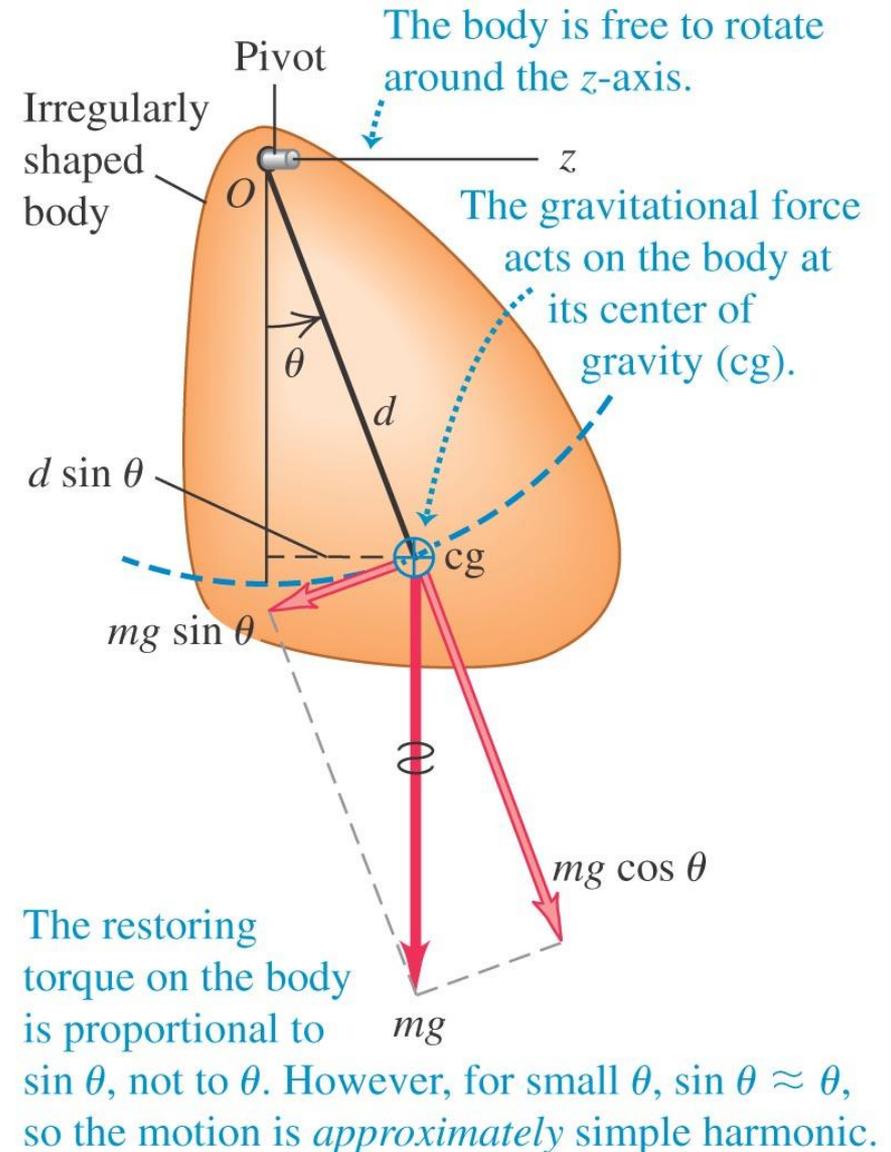
- A *simple pendulum* consists of a point mass (the bob) suspended by a massless, unstretchable string.
- If the pendulum swings with a small amplitude  $\theta$  with the vertical, its motion is simple harmonic.
- $I_{\text{cm}} = mL^2$ ,  $I$  = moment inertia =  $mL^2$
- $\tau = L * m * g \sin(\theta)$
- $\alpha = \text{angular accel} = d^2 \theta / dt^2$
- Eq. motion  $d^2 \theta / dt^2 = (g/L) \sin(\theta) \sim (g/L) \theta$
- Solution is  $\theta(t) = A \sin(\omega t + \phi)$  - SHO
- $A$  – amp,  $\phi$  - phase – both set by initial cond
- $\omega = (g/L)^{1/2}$  angular freq (rad/s)
- $T = 2\pi / \omega = 2\pi (L/g)^{1/2}$
- Note  $T \sim L^{1/2}$  and  $g^{-1/2}$

(b) An idealized simple pendulum



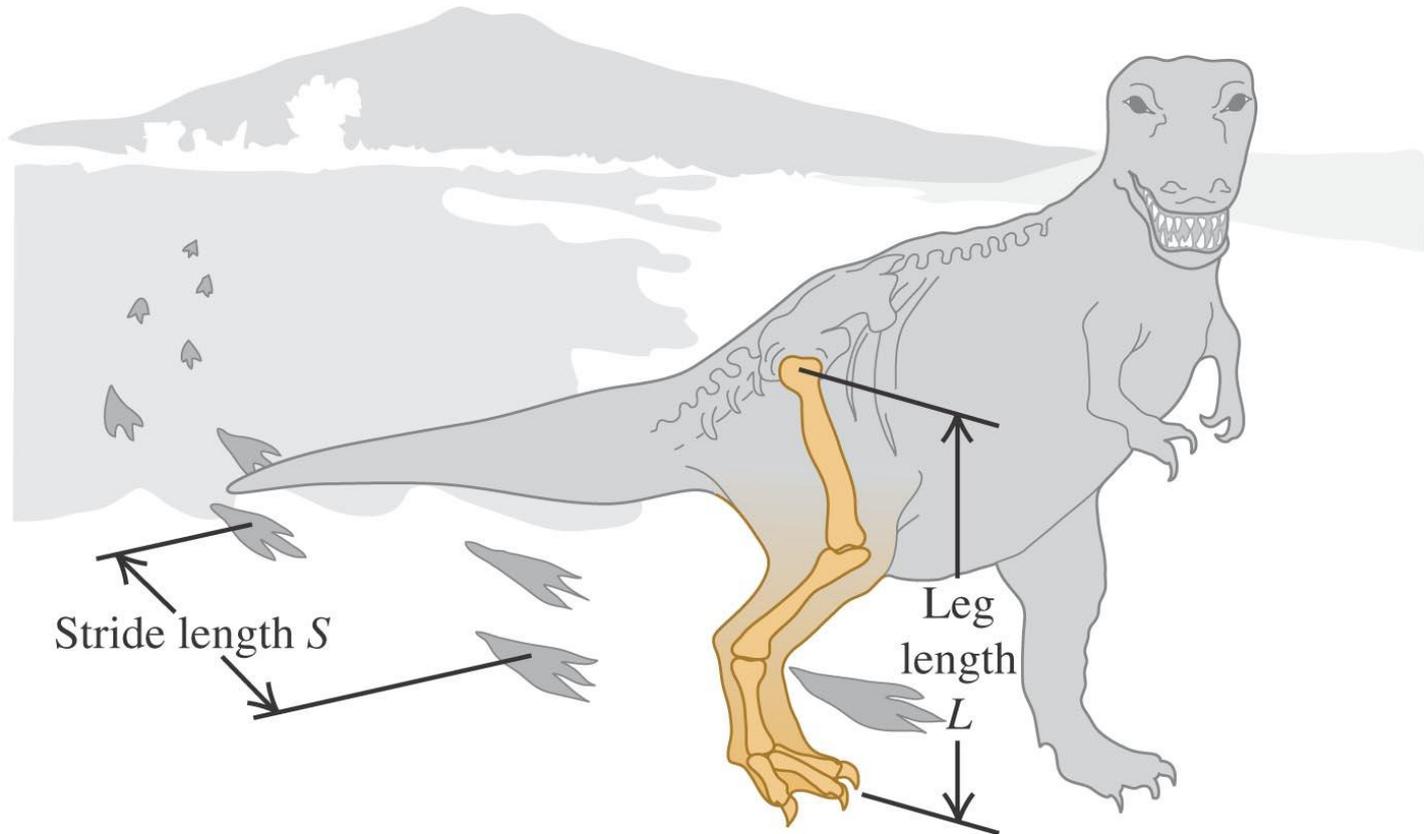
# The physical pendulum

- A *physical pendulum* is any real pendulum that uses an extended body instead of a point-mass bob.
- For small amplitudes, its motion is simple harmonic.
- Same solution as simple pendulum – ie SHO.
- $\omega = (g/L)^{1/2}$  angular freq (rad/s)
- $T = 2\pi/\omega = 2\pi (L/g)^{1/2}$



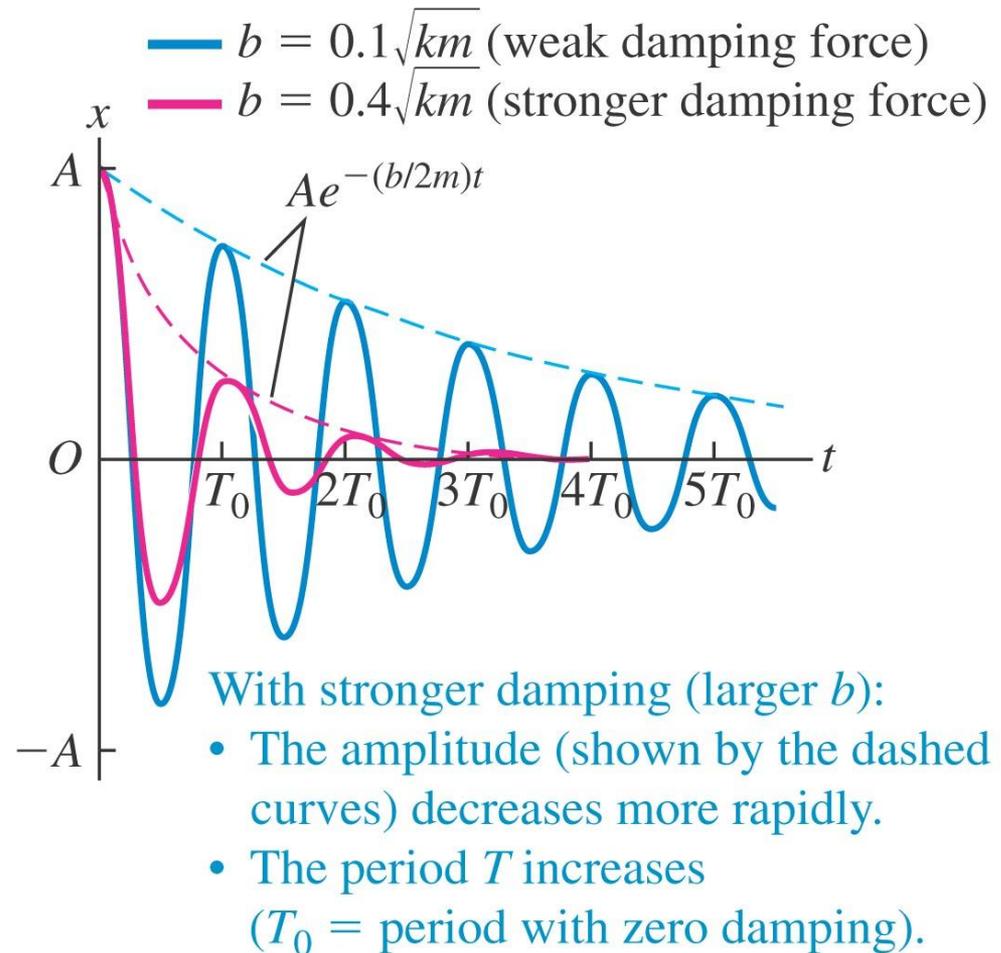
# *Tyrannosaurus rex* and the physical pendulum

- We can model the leg of *Tyrannosaurus rex* as a physical pendulum.
- Unhappy T Rex – cannot use social media in class.



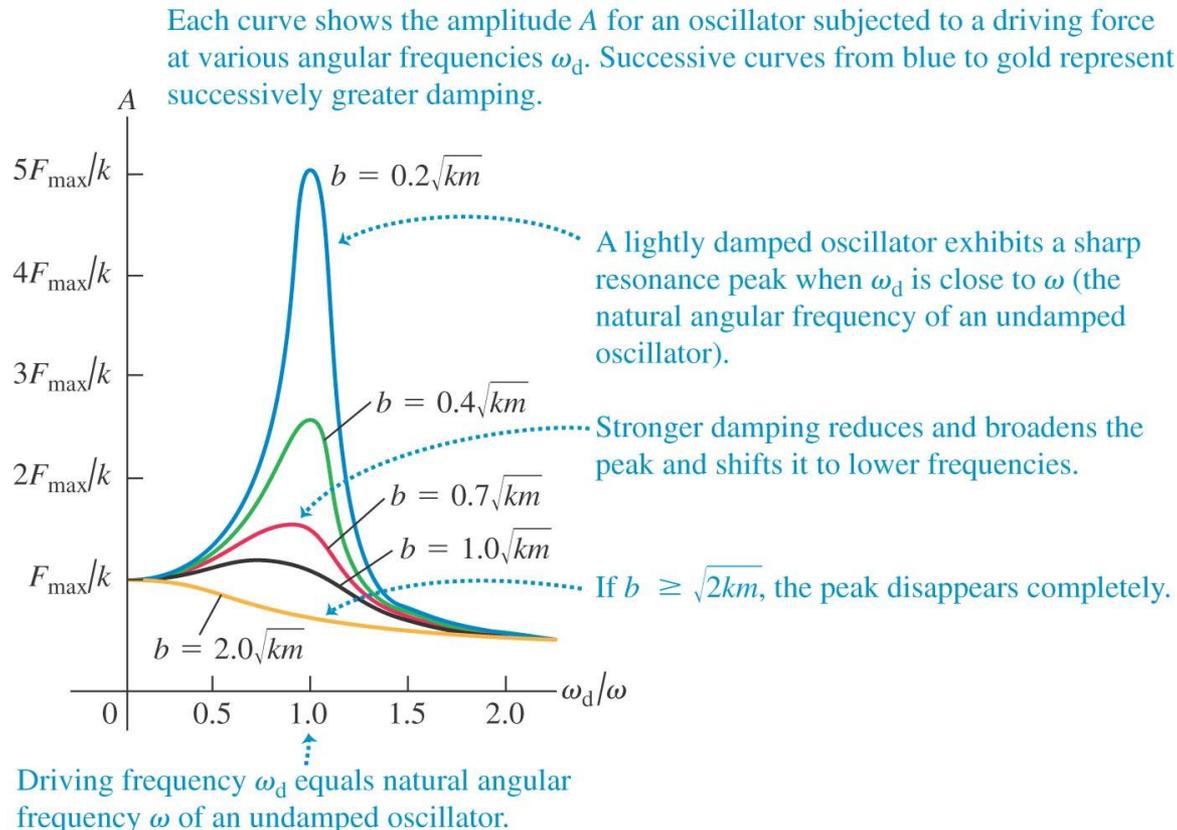
# Damped oscillations

- Real-world systems have some dissipative forces that decrease the amplitude.
- The decrease in amplitude is called *damping* and the motion is called *damped oscillation*.
- Figure illustrates an oscillator with a small amount of damping.
- The mechanical energy of a damped oscillator decreases continuously.



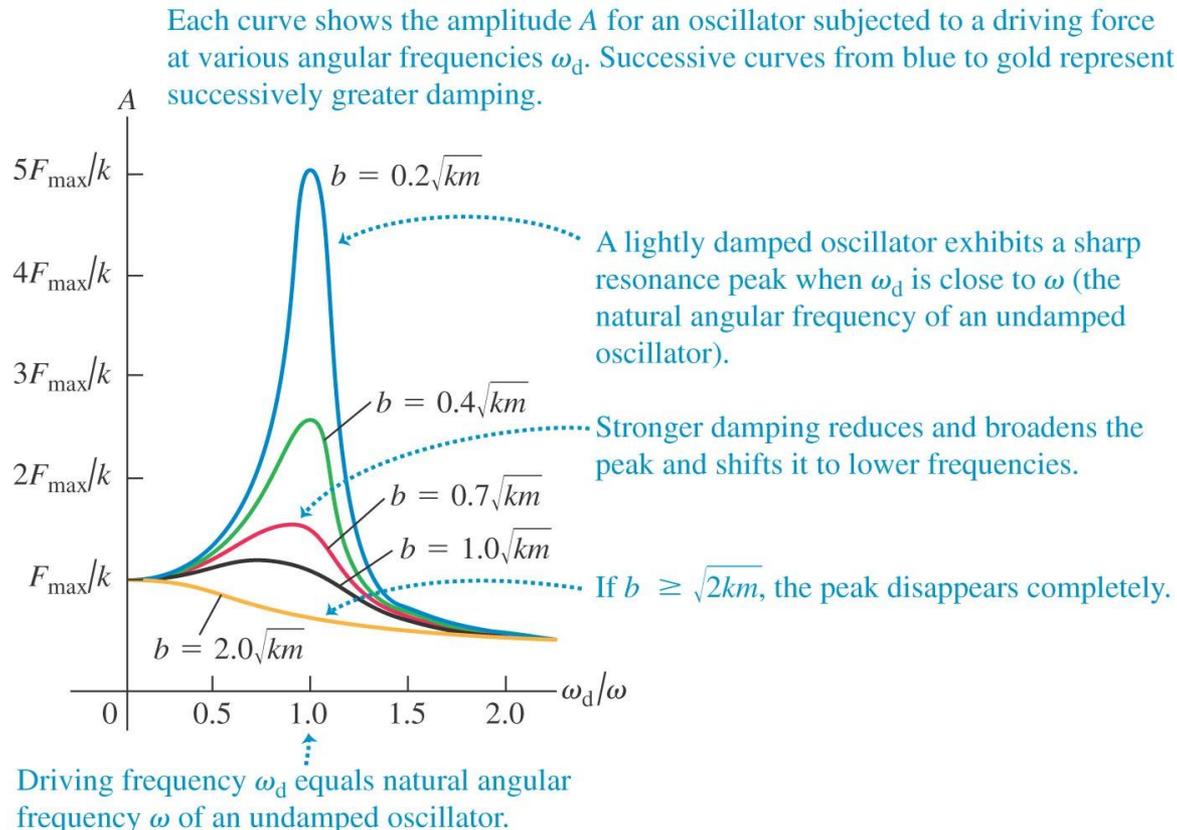
# Forced oscillations and resonance

- A *forced oscillation* occurs if a *driving force* acts on an oscillator.
- *Resonance* occurs if the frequency of the driving force is near the *natural frequency* of the system.

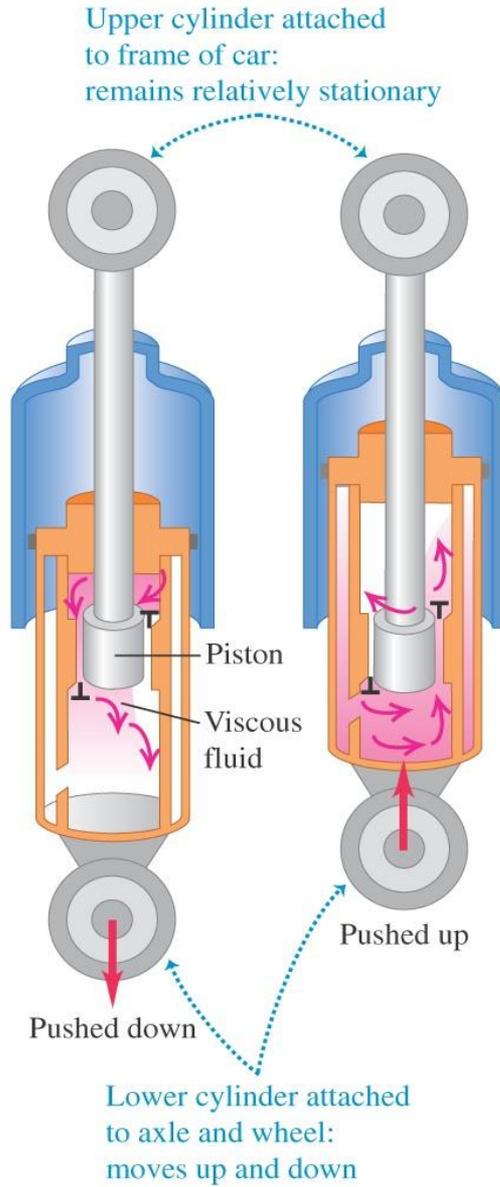


# Forced oscillations and resonance

- A *forced oscillation* occurs if a *driving force* acts on an oscillator.
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# Car shock absorbers - Damped oscillations



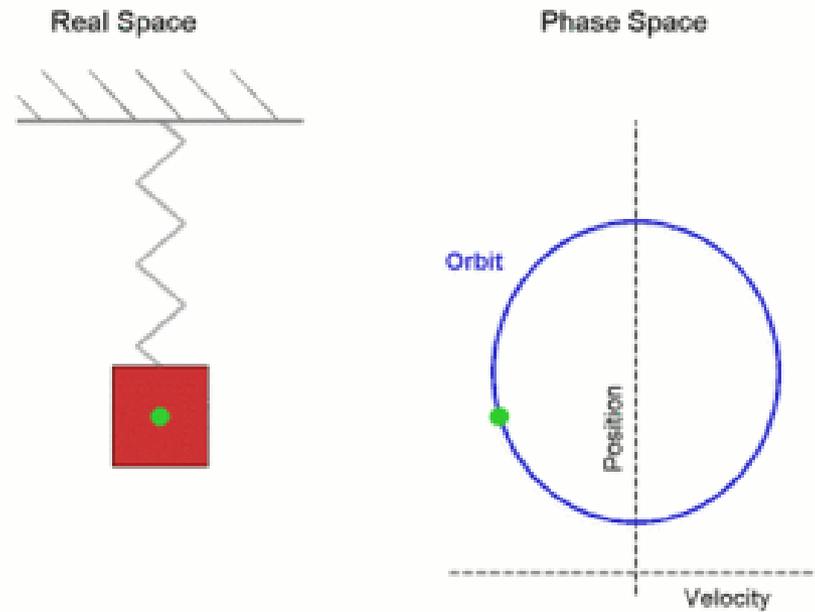
# Forced oscillations and resonance

## Structural Failure

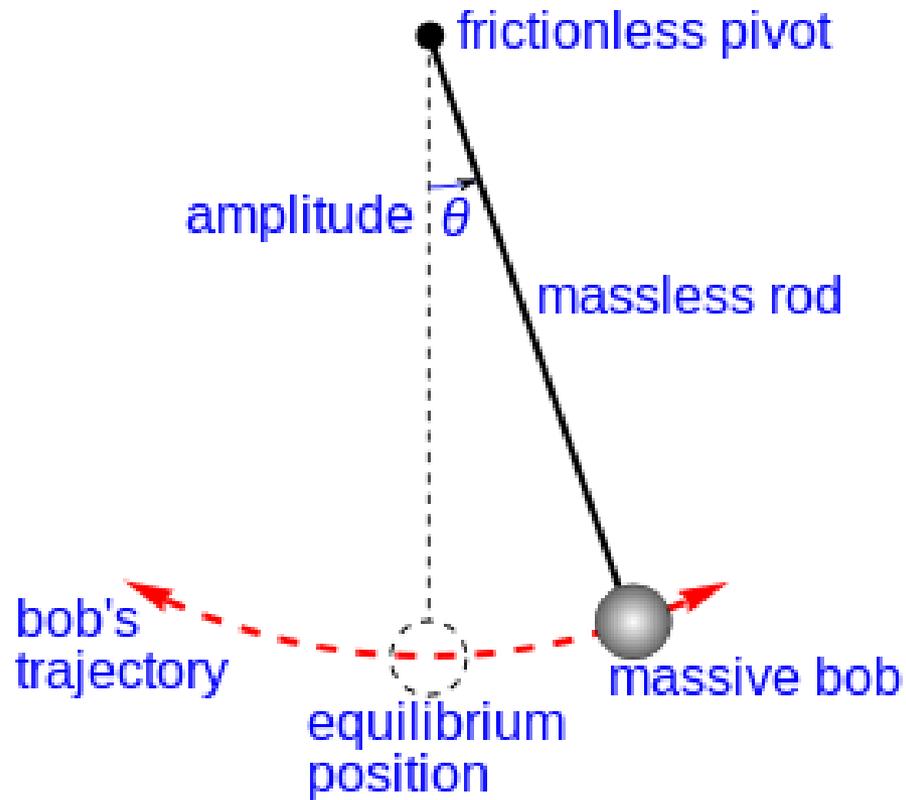
- Nov 7, 1940
- The Tacoma Narrows Bridge suffered spectacular structural failure
- Wind driven osc - too much resonant energy. Too little damping
- <https://www.youtube.com/watch?v=nFzu6CNtqec>



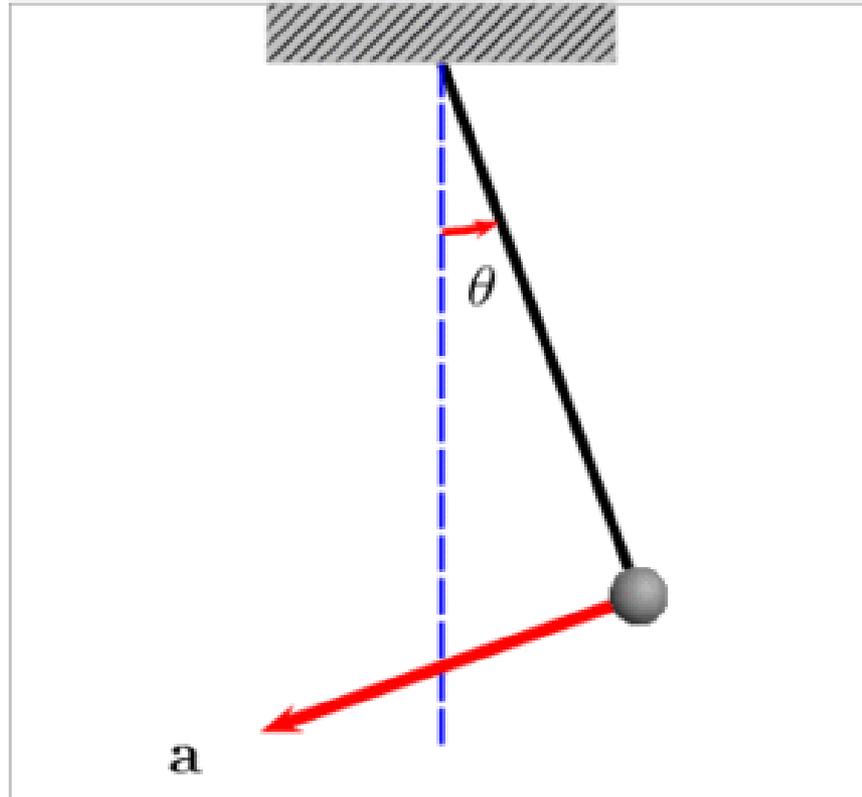
# Simple Harmonic Oscillator (SHO)



# Pendulum

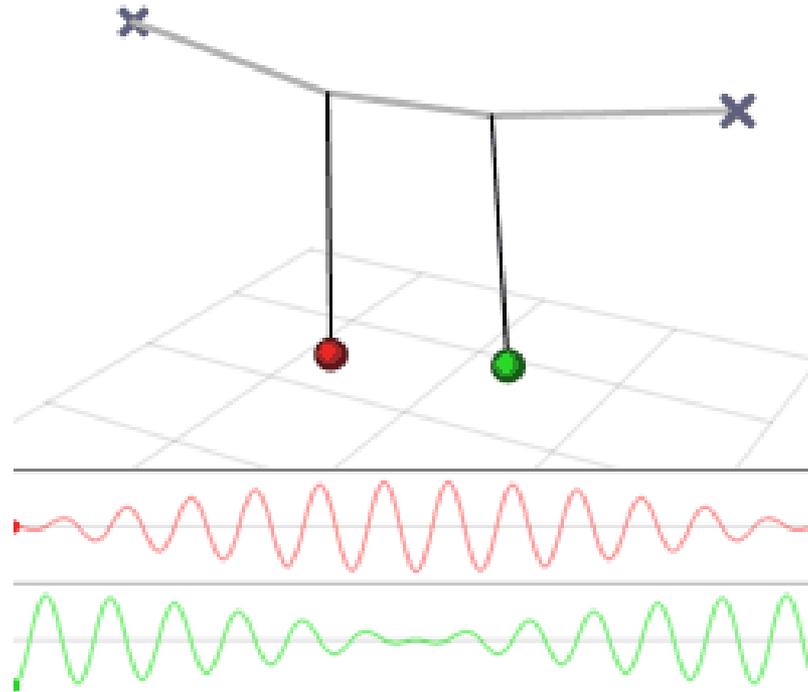


# Simple Pendulum

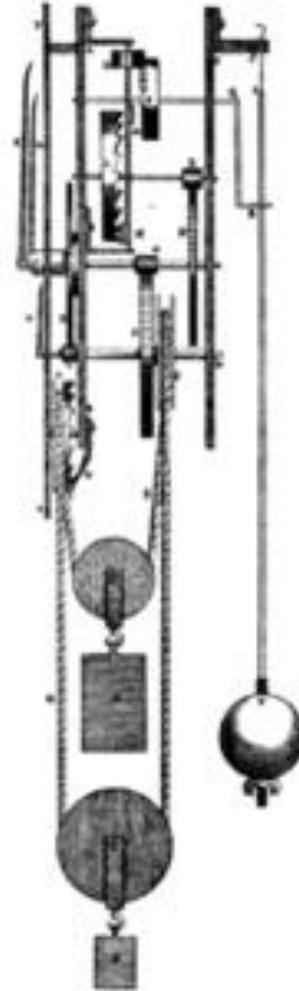
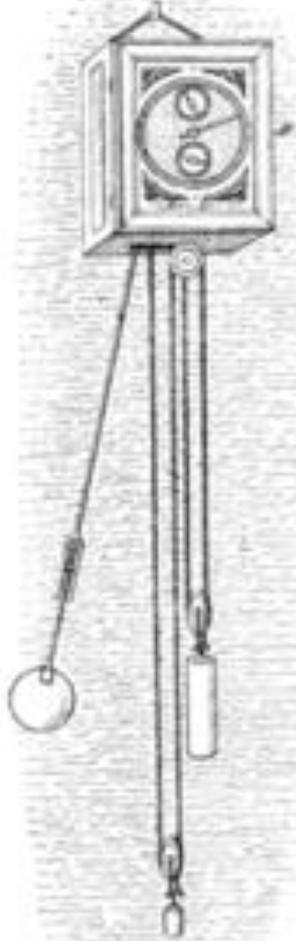


# Two pendulums – same natural freq

## Coupled on wire



# Christian Huygens First Pendulum Clock 1656



# US Time Standard 1909 to 1929

Pendulum is in low pressure vessel

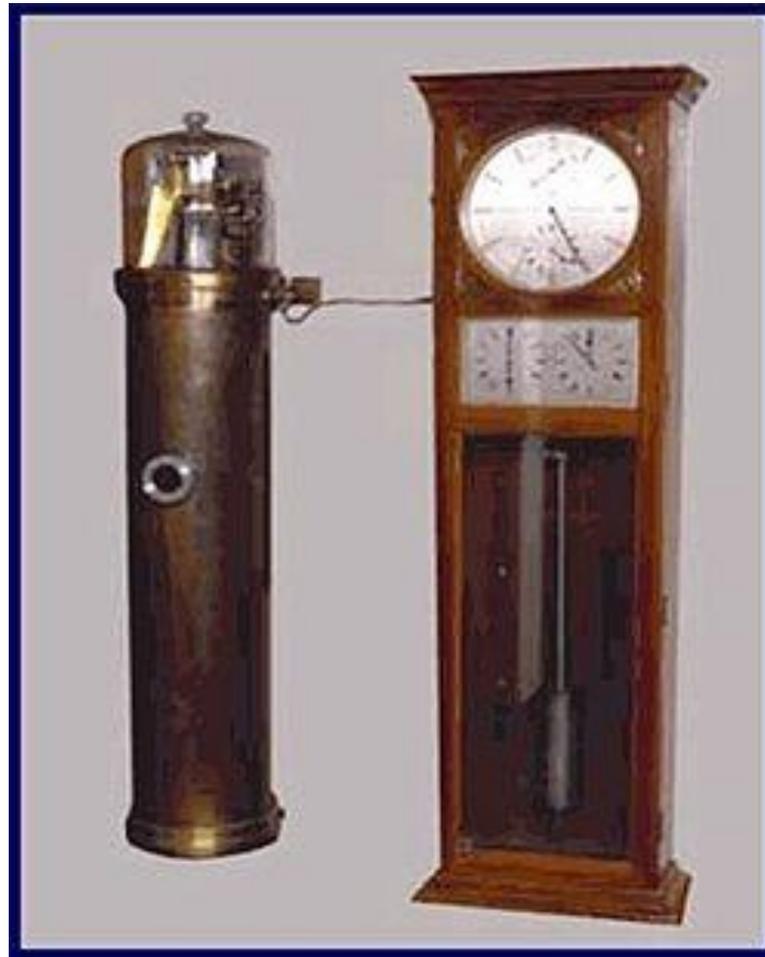
NBS – National Bureau of Standards – now NIST (Natl Inst Sci and Tech)

Riefler regulator



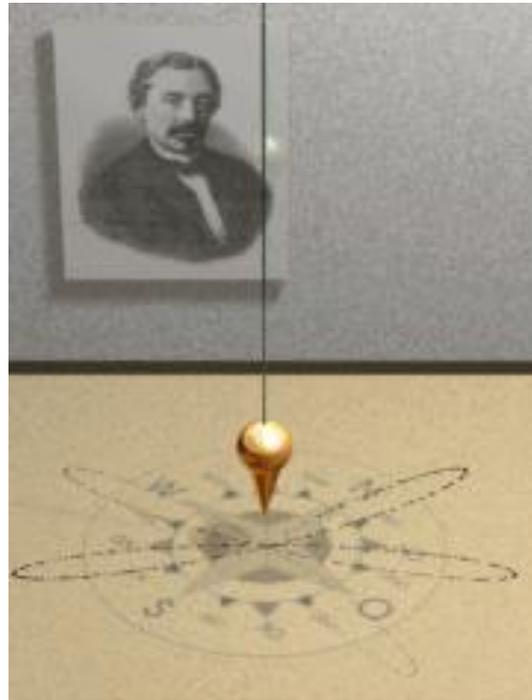
# Vacuum Pendulum – 1 sec / year!!

Synchronized to second pendulum clock



# Foucault Pendulum 1851

Precession of Pendulum  
Showed Earth Rotates



# Seconds Pendulum – 2 sec period Used to Measure Gravity

